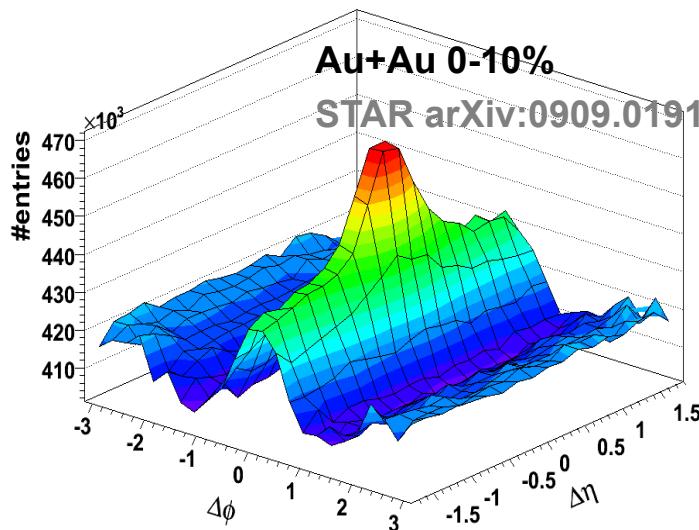
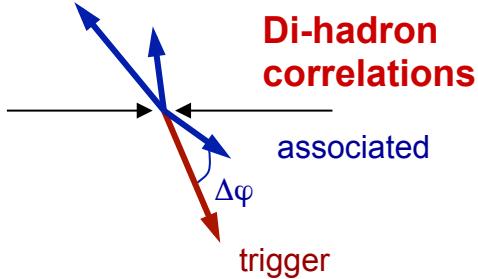


# **Probing extreme QCD through ridge-like correlations in small systems: status and problems**

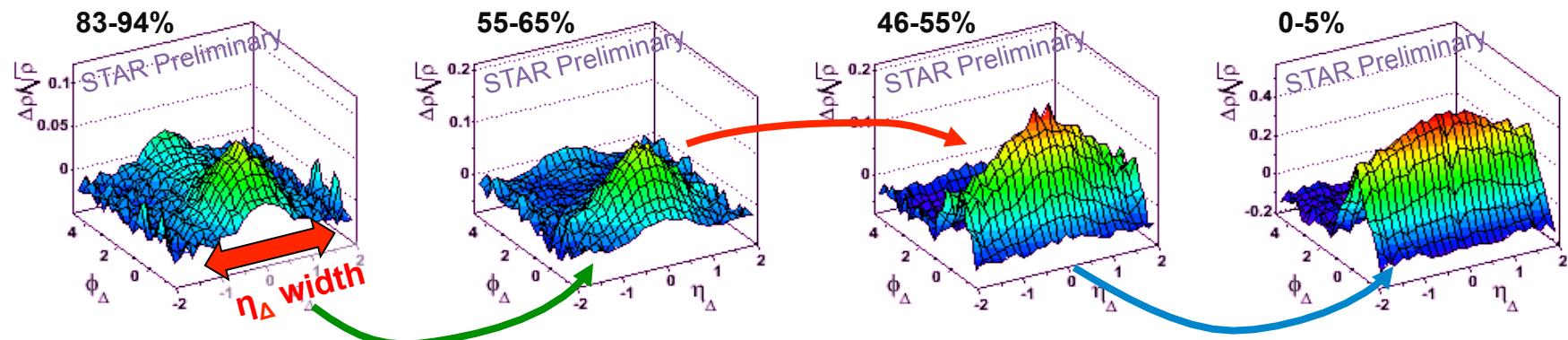
**Raju Venugopalan  
Brookhaven National Laboratory**

**RBRC Saturation Workshop, April 26-28, 2017**

# The ridge in A+A collisions



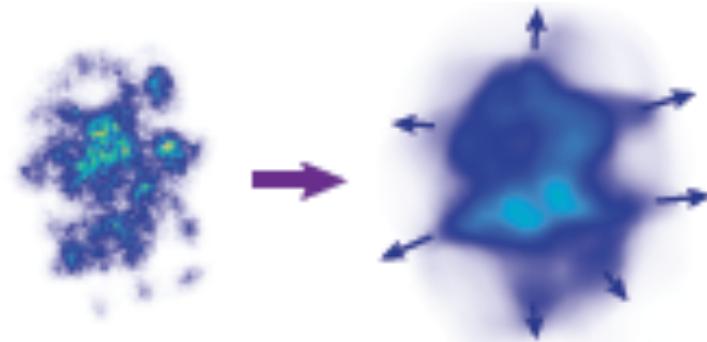
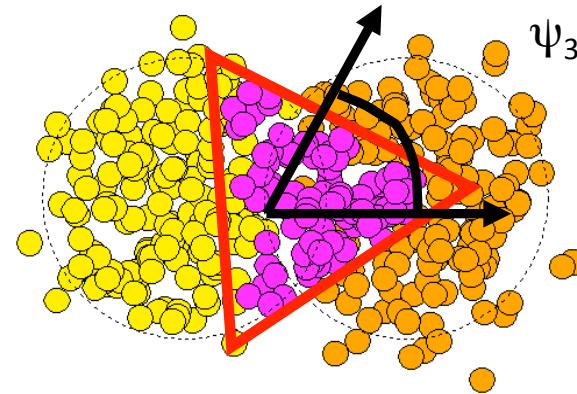
$3 < p_{t,\text{trigger}} < 4 \text{ GeV}$   
 $p_{t,\text{assoc.}} > 2 \text{ GeV}$



Collimated, long range rapidity correlations:  
 First seen by RHIC Au+Au experiments: STAR, PHOBOS, PHENIX

# The ridge in A+A collisions

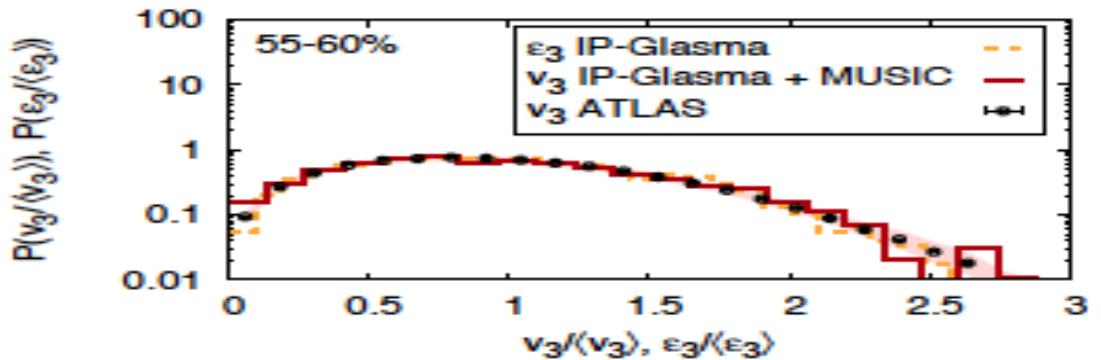
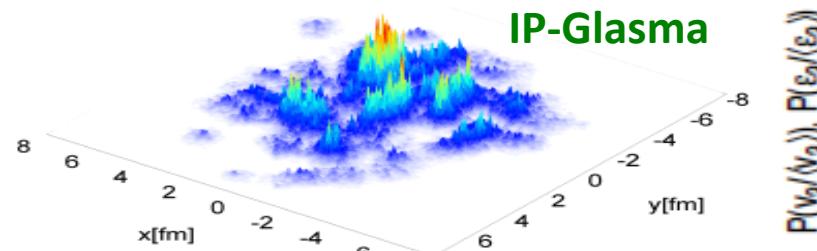
Alver, Roland, PRC81(2010) 054905  
Alver, Gombeaud, Luzum, Ollitrault, PRC82 (2010) 03491



Structure of ridge-correlations can be understood as hydrodynamic flow driven by event-by-event fluctuations in nucleon positions

$$\frac{1}{N_{\text{trig}} N_{\text{assoc}}} \frac{d^2 N}{d\Delta\Phi} = 1 + V_1 \cos(\Delta\Phi) + V_2 \cos(2\Delta\Phi) + \dots$$

Gale, Jeon, Schenke, Tribedy, Venugopalan, PRL110 (2013) 012302



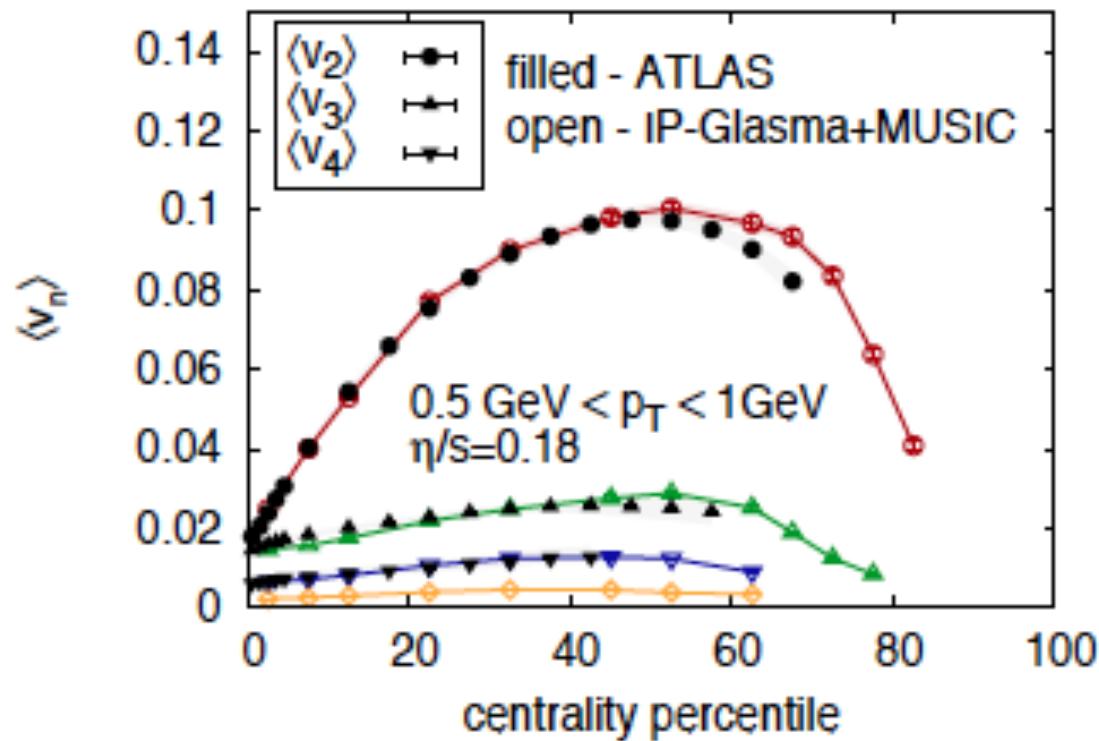
Some evidence of sensitivity of data to sub-nucleon scale fluctuations

# What's the smallest sized QGP droplet?

IP-Glasma= initial state

MUSIC=event.by.event. hydro

Schenke, Venugopalan, PRL 113 (2014) 102301



Where does the hydro paradigm break down?

# Higher cumulants of elliptic flow

m-particle flow  
cumulants

$$c_n \{2m\} = \langle\langle e^{in(\phi_1 + \cdots + \phi_m - \phi_{m+1} - \cdots - \phi_{2m})} \rangle\rangle$$

Borghini,Dinh,Ollitrault, nucl-th/0105040

$$v_n \{2\}^2 \equiv c_n \{2\}; v_n \{4\}^2 \equiv -c_n \{4\}; v_n \{6\}^6 \equiv c_n \{6\}/4$$

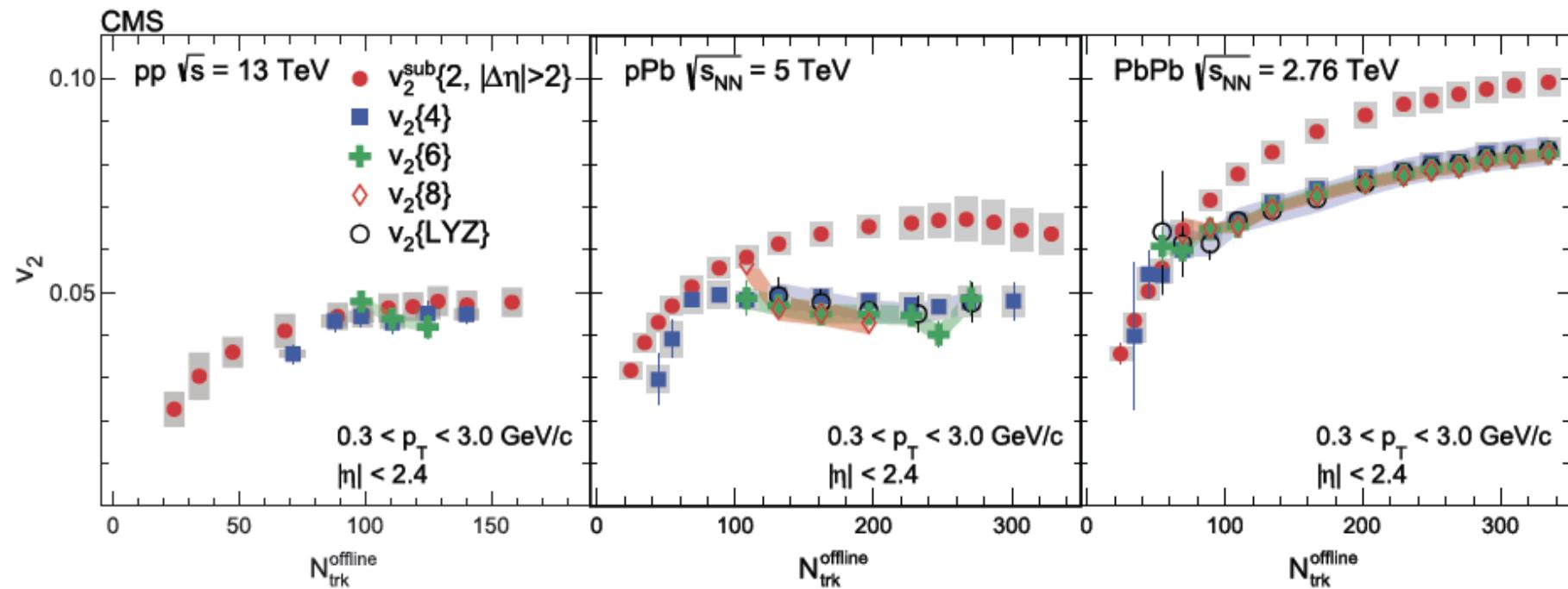
Spatial eccentricities:  $\epsilon_n = \frac{1}{\langle r_\perp^n \rangle} \int d^2 r_\perp e^{in\phi_r} r_\perp^n \frac{dN}{dy d^2 r_\perp}$

A number of simple models give  $\epsilon_n \{2\} > \epsilon \{4\} = \epsilon \{6\} = \dots$

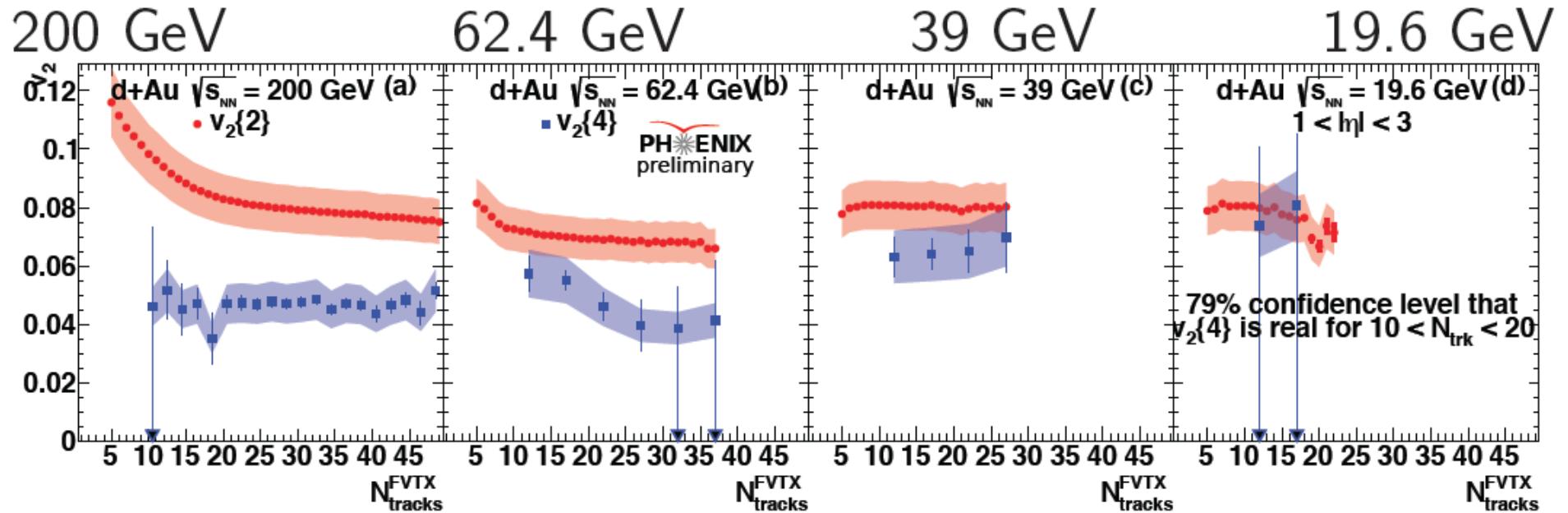
Hydro linear response:  $v_n \{m\} \approx c_n \epsilon_n \{m\}$

Gardim,Grassi,Luzum,Ollitrault, PRC (2012)024908; Niemi,Denicol,Holopainen,Huovinen, PRC87 (2013)054901  
Bzdak,Bozek,McLerran, arXiv:1311.7325, Bzdak, Skokov, arXiv: 1312.7349  
Li, Ollitrault, arXiv:1312.6555, Basar,Teaney, arXiv:1312.6770

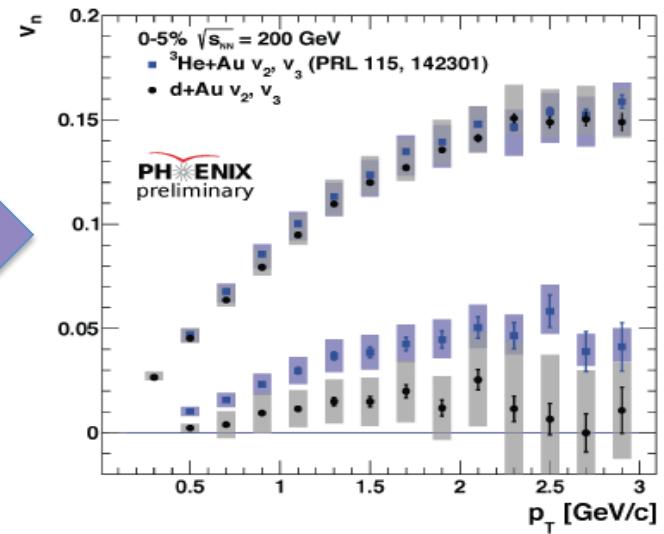
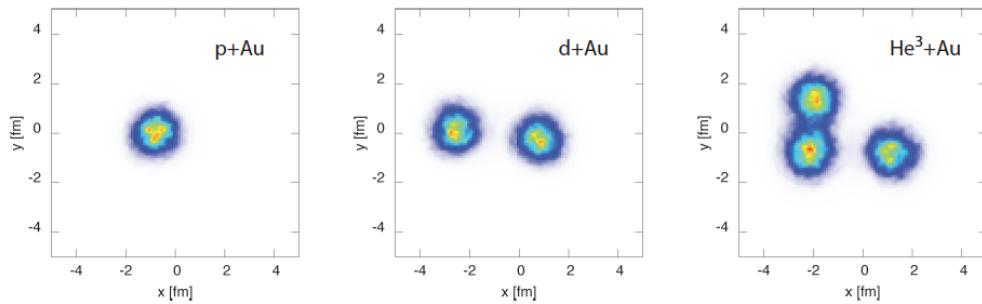
# Collectivity across system size



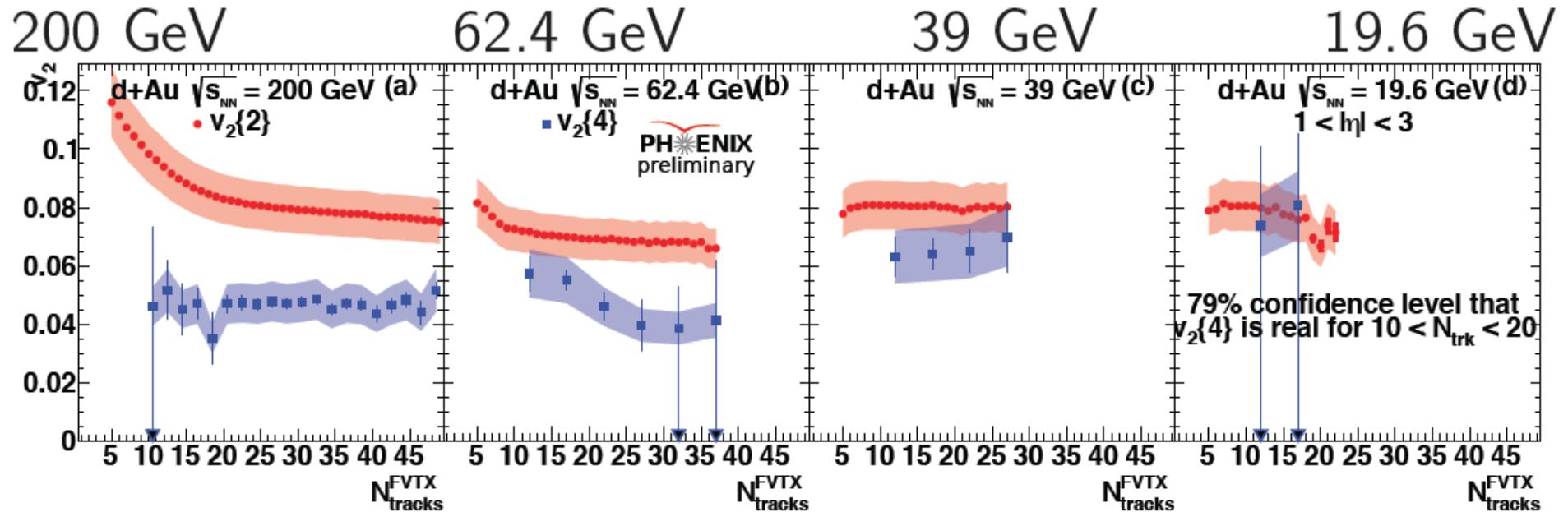
# Collectivity across wide energy scales



Schenke, RV:1407.7557



# Collectivity across wide energy scales



Panta Rhei ?



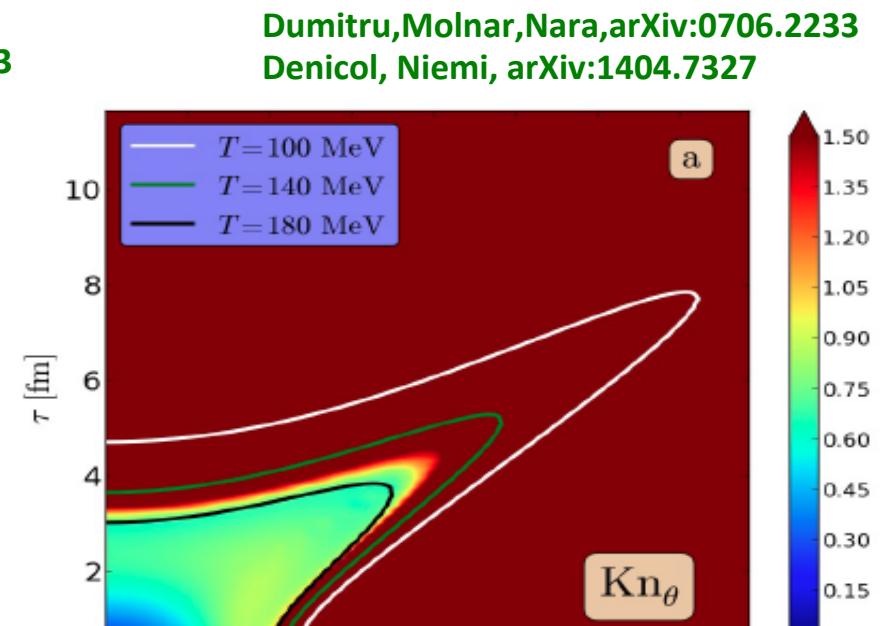
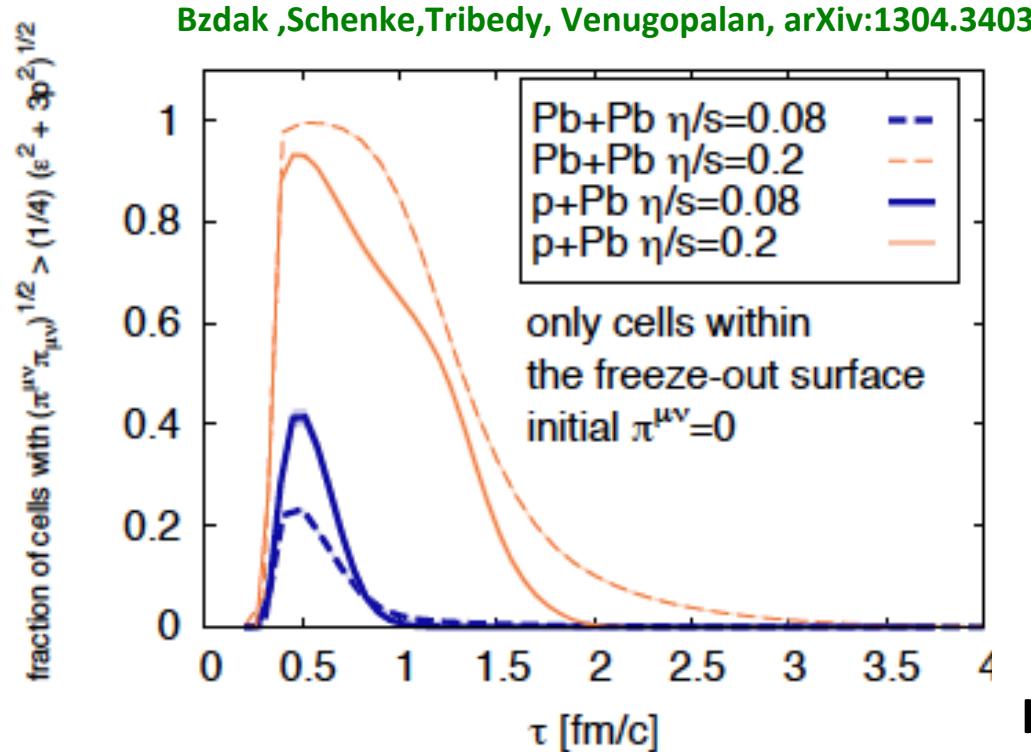
Heraclitus of Ephesus  
535-475 BC

# Issues with the hydrodynamic paradigm: I

Two frequently used measures: Reynolds # and Knudsen #

$$R^{-1} \propto (\Pi^{\mu\nu}\Pi_{\mu\nu})^{1/2}/(\epsilon^2 + 3P^2)^{1/2}$$

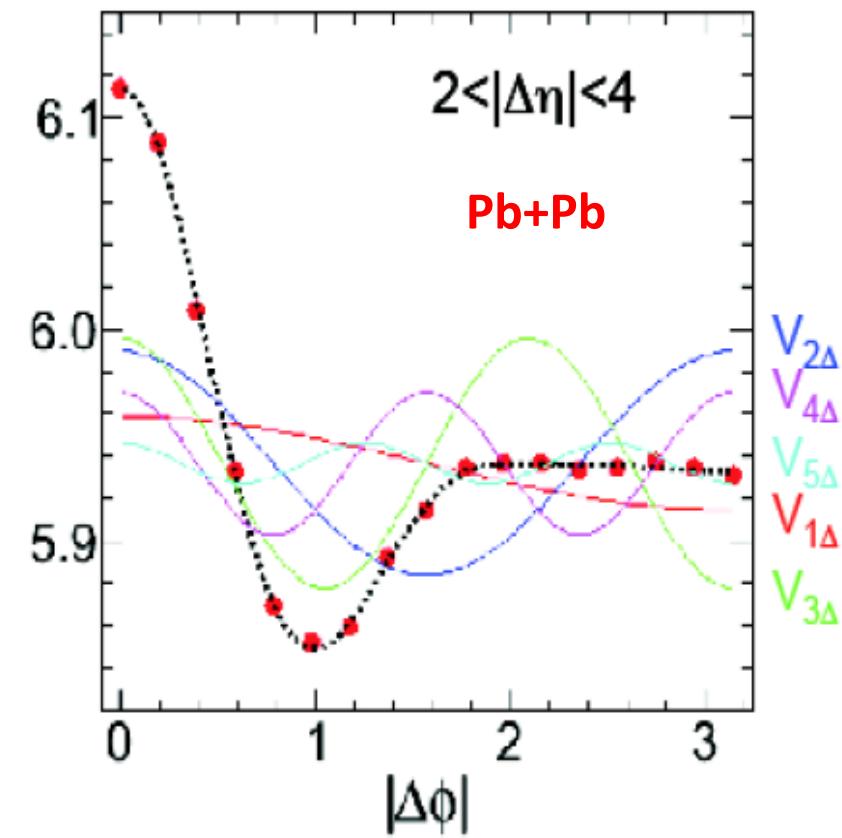
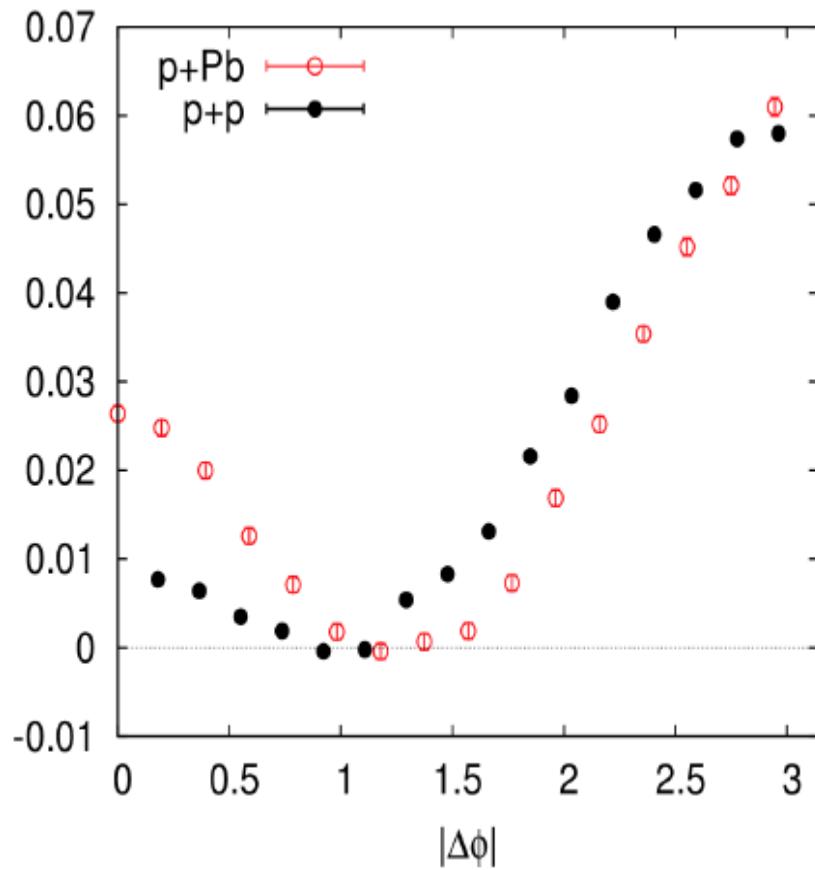
$$\text{Kn} = \frac{\tau_\pi}{L} ; \quad \tau_\pi \propto \frac{\eta}{sT}$$



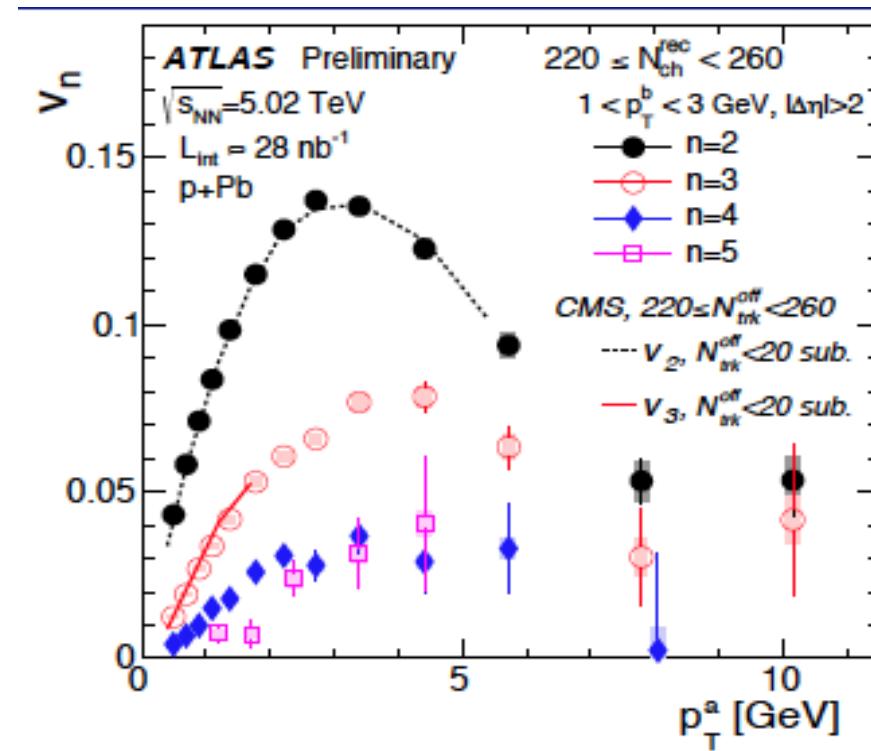
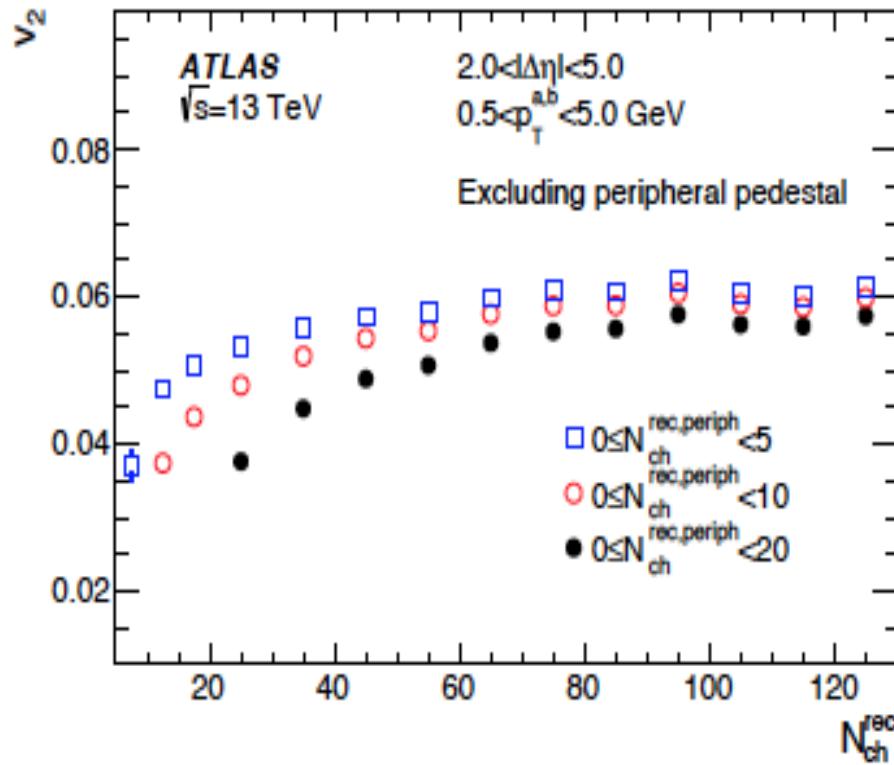
Hydro good for  $\text{Kn} < 0.5$ ,  
marginal for  $\text{K} < 1$  transient regime;  
 $\text{K} > 1$  free streaming

## Issues with the hydrodynamic paradigm: II

No (mini-) jet quenching seen in the smaller systems



# Issues with the hydrodynamic paradigm: III



Large anisotropies at larger  $p_T$  and smaller  $N_{\text{ch}}$  than one might reconcile with a hydrodynamic description

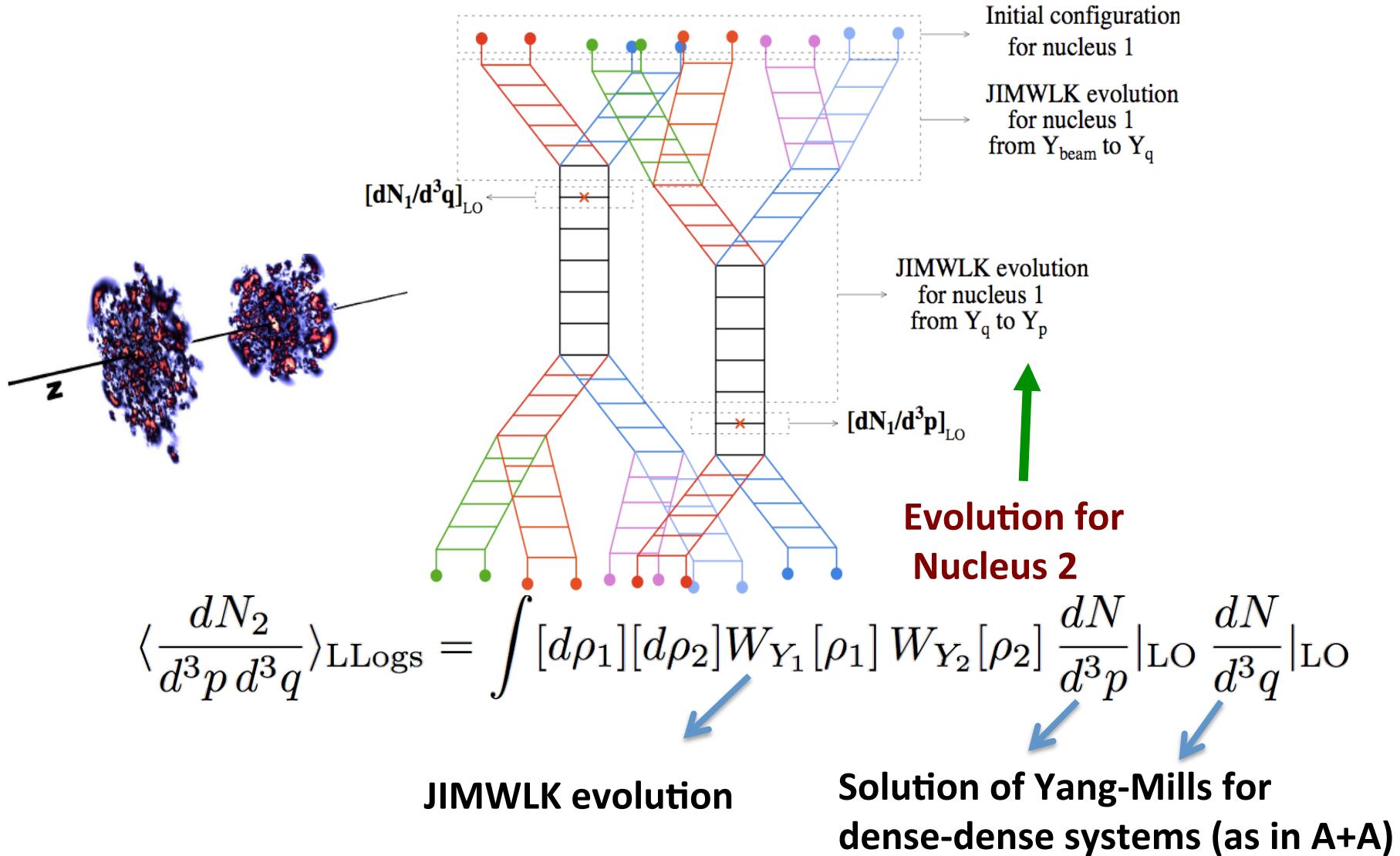
**Can we understand multiparticle correlations in  
an *ab initio* approach**

# Two-parton azimuthal correlations in the CGC

Dumitru,Gelis,McLerran,Venugopalan: 0804.3858

Gelis, Lappi, Venugopalan, arXiv: 0807.1306

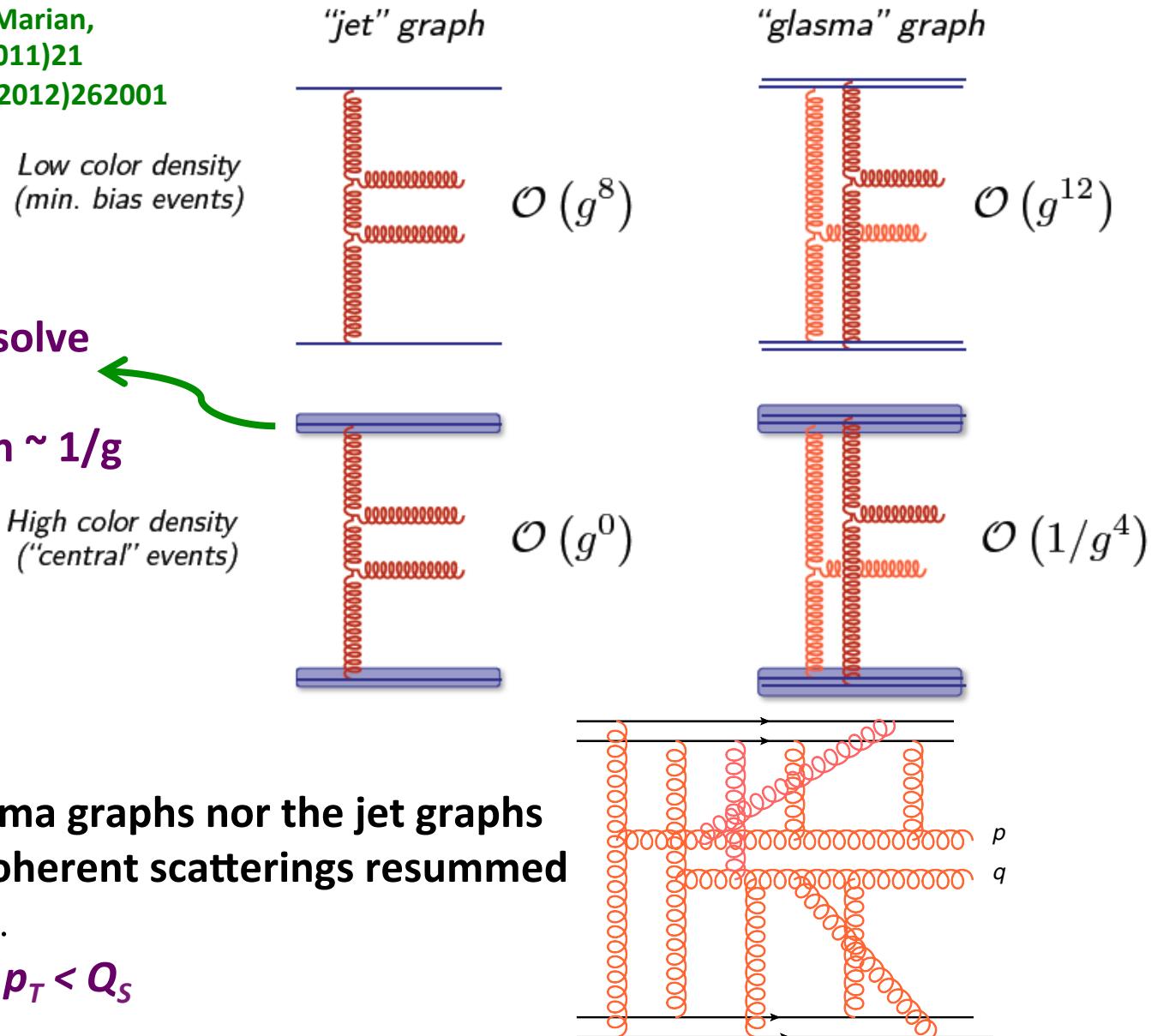
Dusling,Gelis,Lappi,Venugopalan, arXiv:0911.2720



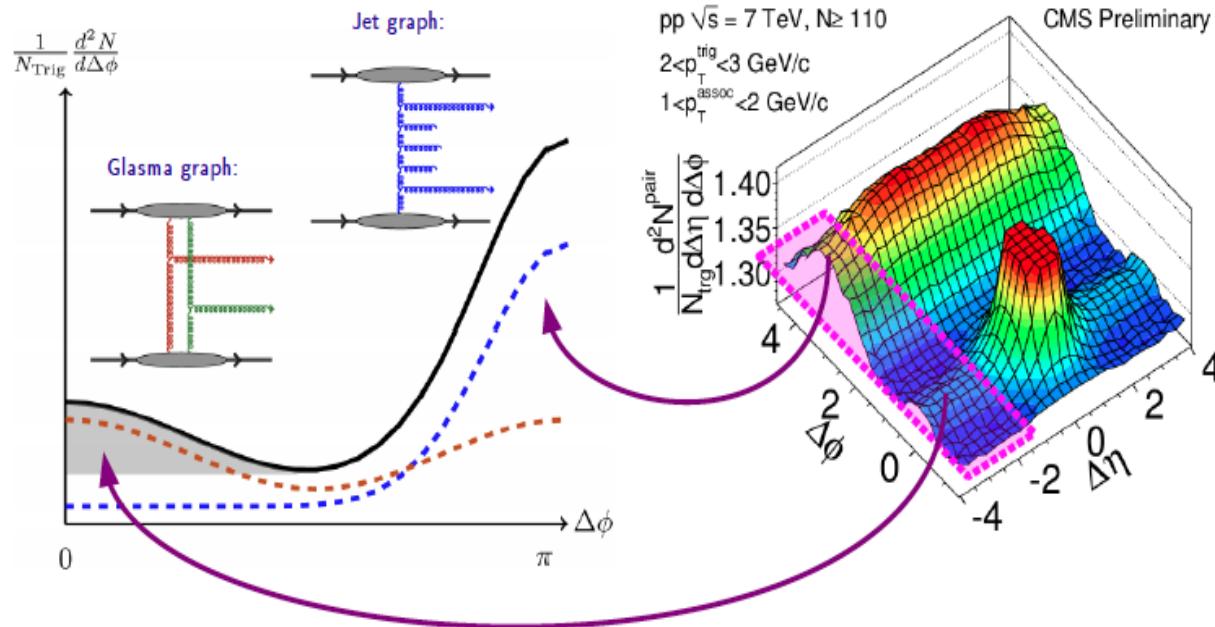
# Glasma graph approximation: power counting

Dumitru,Dusling,Gelis,Jalilian-Marian,  
Lappi,Venugopalan, PLB697 (2011)21  
Dusling,Venugopalan,PRL108 (2012)262001

Gluons with  $k_T \sim Q_S$  resolve  
 $n \sim 1/g^2$  color sources  
Effective coupling:  $g^*n \sim 1/g$



# Anatomy of long range collimations



**RG evolution of Glasma graphs:**

$$C(\mathbf{p}, \mathbf{q}) \propto \frac{g^4}{\mathbf{p}_\perp^2 \mathbf{q}_\perp^2} \int d^2 \mathbf{k}_{1\perp} \Phi_{A_1}^2(y_p, \mathbf{k}_{1\perp}) \Phi_{A_2}(y_p, \mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_2}(y_q, \mathbf{q}_\perp - \mathbf{k}_{1\perp})$$

+ permutations

Dusling, RV, PRD 87, 051502 (R) (2013); PRD87 (2013) 094034  
 Dusling, Tribedy, RV, PRD93 (2016) 014034

**RG evolution of the mini-jets:**

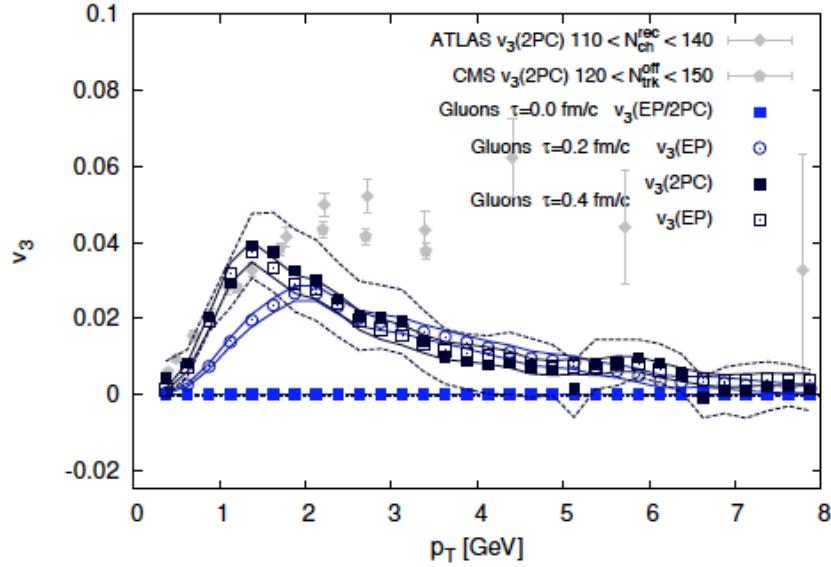
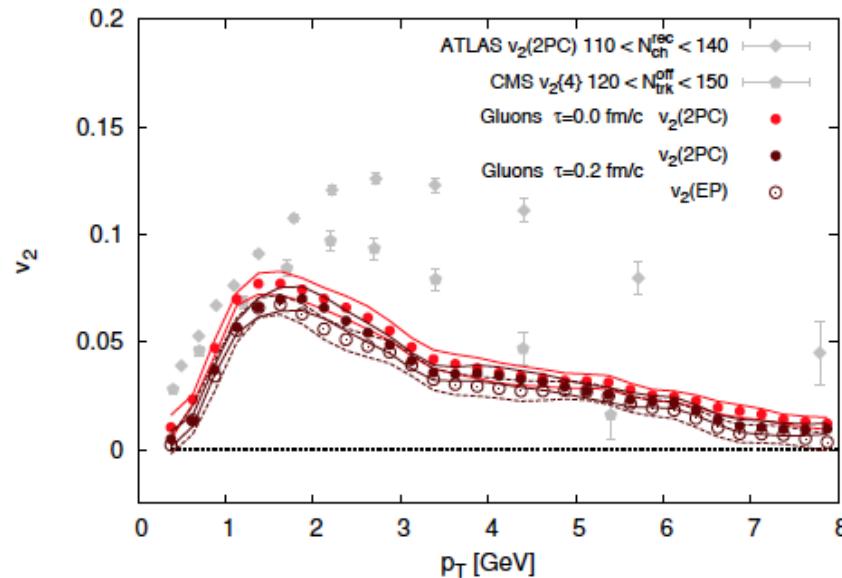
$$C_{\text{dijet}}(\mathbf{p}, \mathbf{q}) \propto \Phi_A \otimes \Phi_B \otimes G_{\text{BFKL}}$$

**Good agreement with data for  $p_T > Q_S$**

**However no odd harmonics  $v_3, v_5$  in this approximation...**

# Azimuthal anisotropy from Yang-Mills dynamics

Schenke,Schlichting,RV, PLB747(2015)76



See also, Lappi,Srednyak,RV, JHEP1001 (2010 )066

Recent analytical work in dilute-dense approx:

Kovchegov,Wertepny,NPA906 (2013)50

McLerran, Skokov arXiv:1611.09870

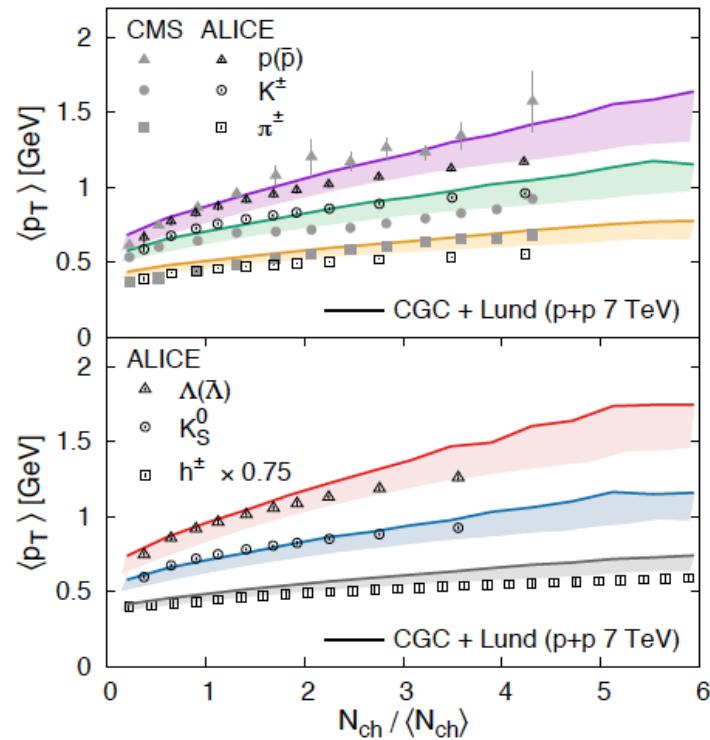
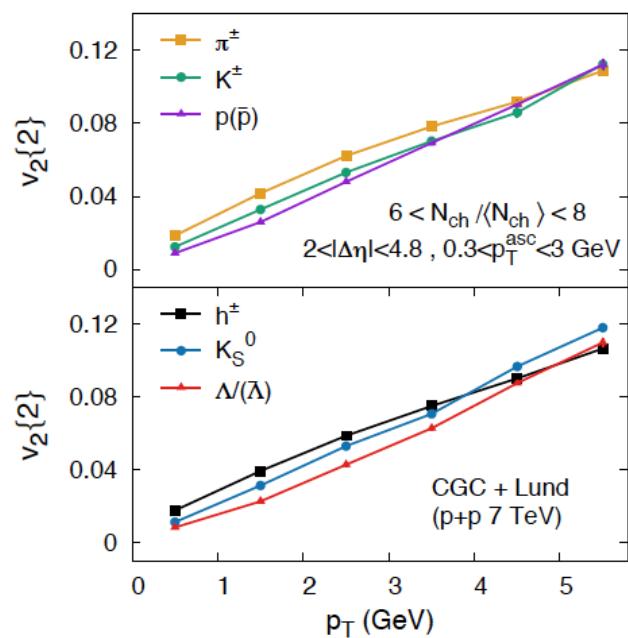
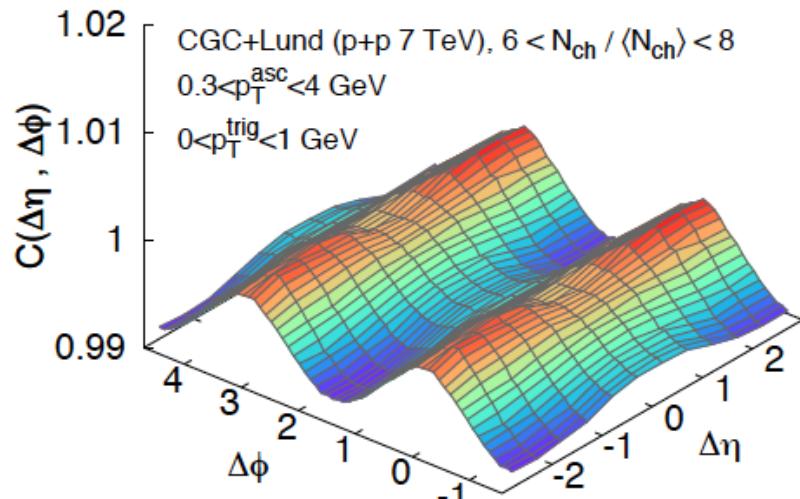
Kovner,Lublinsky,Skokov, arXiv:1612.07790

What about 4 and higher particle cumulants from YM dynamics?

Numerically very challenging to look at rare events – in progress

Schenke,Schlichting,Tribedy,RV

# IP-Glasma+Lund fragmentation



Schenke, Schlichting, Tribedy, RV, PRL117(2016)162301

**Pattern of mass splitting of  $\langle p_T \rangle$  and  $v_2$  seen in high multiplicity events is reproduced**

# Tracing azimuthal initial state correlations

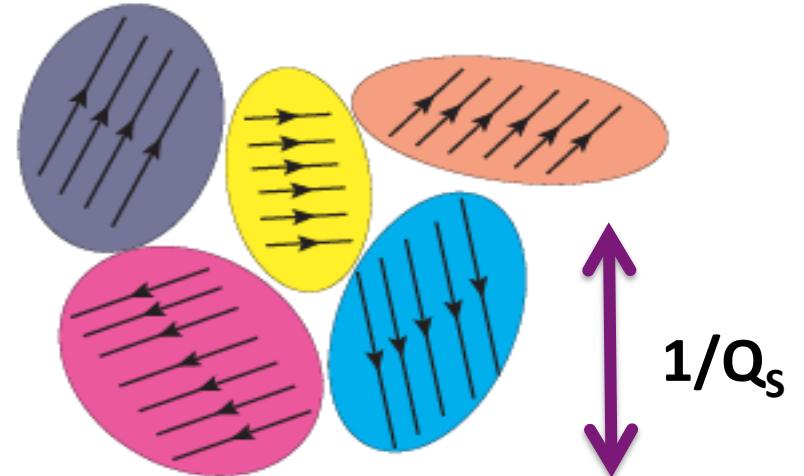
Multi-particle correlations from eikonal scattering of partons off color domains in a nuclear target

Kovner,Lublinsky,arXiv:1012.3398,1109.0347

Dumitru,Gianini, arXiv:1406.5781

Dumitru,Skokov,arXiv:1411.6630,

Dumitru,McLerran,Skokov,arXiv:1410.4844



Explore in simple model:

$$\frac{d^2 N}{d^2 \mathbf{p}_1 d^2 \mathbf{p}_2} = \int d^2 \mathbf{b}_1 d^2 \mathbf{b}_2 \int \frac{d^2 \mathbf{k}_1}{(2\pi)^2} \frac{d^2 \mathbf{k}_2}{(2\pi)^2} W(\mathbf{b}_1, \mathbf{k}_1, \mathbf{b}_2, \mathbf{k}_2) \\ \times \int d^2 \mathbf{r}_1 d^2 \mathbf{r}_2 e^{i(\mathbf{p}_1 - \mathbf{k}_1)} e^{i(\mathbf{p}_2 - \mathbf{k}_2)} \langle D\left(\mathbf{b}_1 + \frac{\mathbf{r}_1}{2}, \mathbf{b}_1 - \frac{\mathbf{r}_1}{2}\right) D\left(\mathbf{b}_2 + \frac{\mathbf{r}_2}{2}, \mathbf{b}_2 - \frac{\mathbf{r}_2}{2}\right) \rangle$$

---

**Dipole**  $D(\mathbf{x}, \mathbf{y}) = \frac{1}{N_c} \text{Tr} (V(\mathbf{x}) V^\dagger(\mathbf{y})) \quad V(\mathbf{x}) = P \exp \left( -ig \int dx^+ A^-(\mathbf{x}, x^+) \right)$

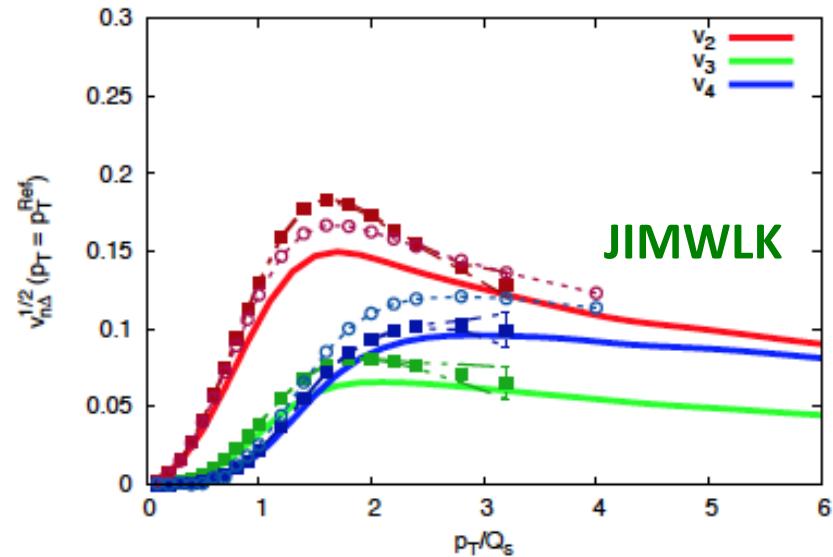
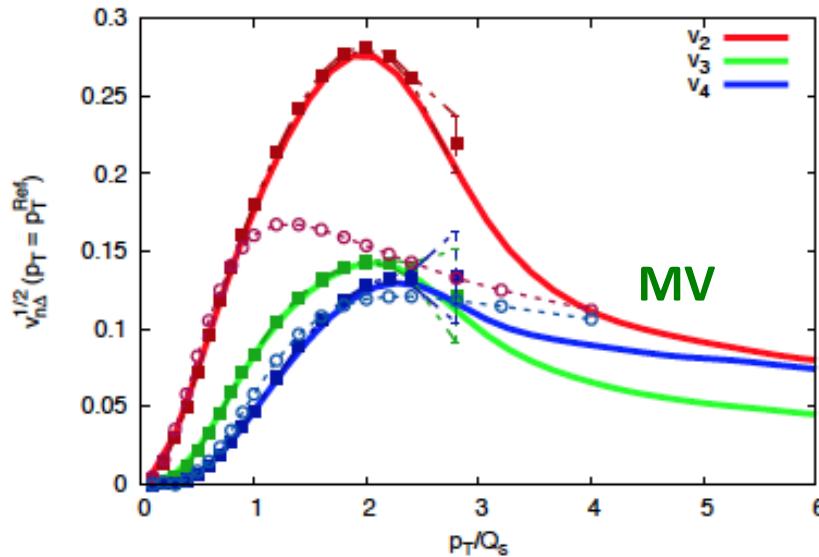
**Ansatz:**  $W(\mathbf{b}_1, \mathbf{k}_1, \mathbf{b}_2, \mathbf{k}_2) = W(\mathbf{b}_1, \mathbf{k}_1) W(\mathbf{b}_2, \mathbf{k}_2)$

$W(\mathbf{b}, \mathbf{k}) = \exp(-\mathbf{b}^2/B - \mathbf{k}^2 B)$  **B=transverse area of projectile**

# Tracing azimuthal initial state correlations

Lappi, arXiv:1501.05505

Lappi,Schenke,Schlichting,RV, arXiv:1509.03499



What about 4-particle correlations?

$$\frac{d^m N}{\prod_{i=1}^m d^2 p_{i\perp}} = \frac{1}{(4\pi^3 B)^m} \prod_{i=1}^m \int d^2 b_i \int d^2 r_i e^{-b_i^2/E} e^{-r_i^2/4B} e^{i p_{i\perp} \cdot r_i} \left\langle \prod_{j=1}^m D(b_j + \frac{r_j}{2}, b_j - \frac{r_j}{2}) \right\rangle$$

$$\kappa_n\{m\} = \int d^2 p_{1\perp} \dots d^2 p_{m\perp} \frac{d^m N}{\prod_{k=1}^m d^2 p_{k\perp}} \prod_{j=1}^{m/2} \prod_{l=\frac{m}{2}+1}^m e^{ni(\phi_j^p - \phi_l^p)}$$

$$c_n\{4\} = \frac{\kappa_n\{4\}}{\kappa_0\{4\}} - 2 \left( \frac{\kappa_n\{2\}}{\kappa_0\{2\}} \right)^2$$

First initial state results on  $c_n\{4\}$ , SC{m,n}  
Kevin Dusling's talk at 10:30 am Friday  
Dusling,Mace,RV, in preparation

## **Summary-I**

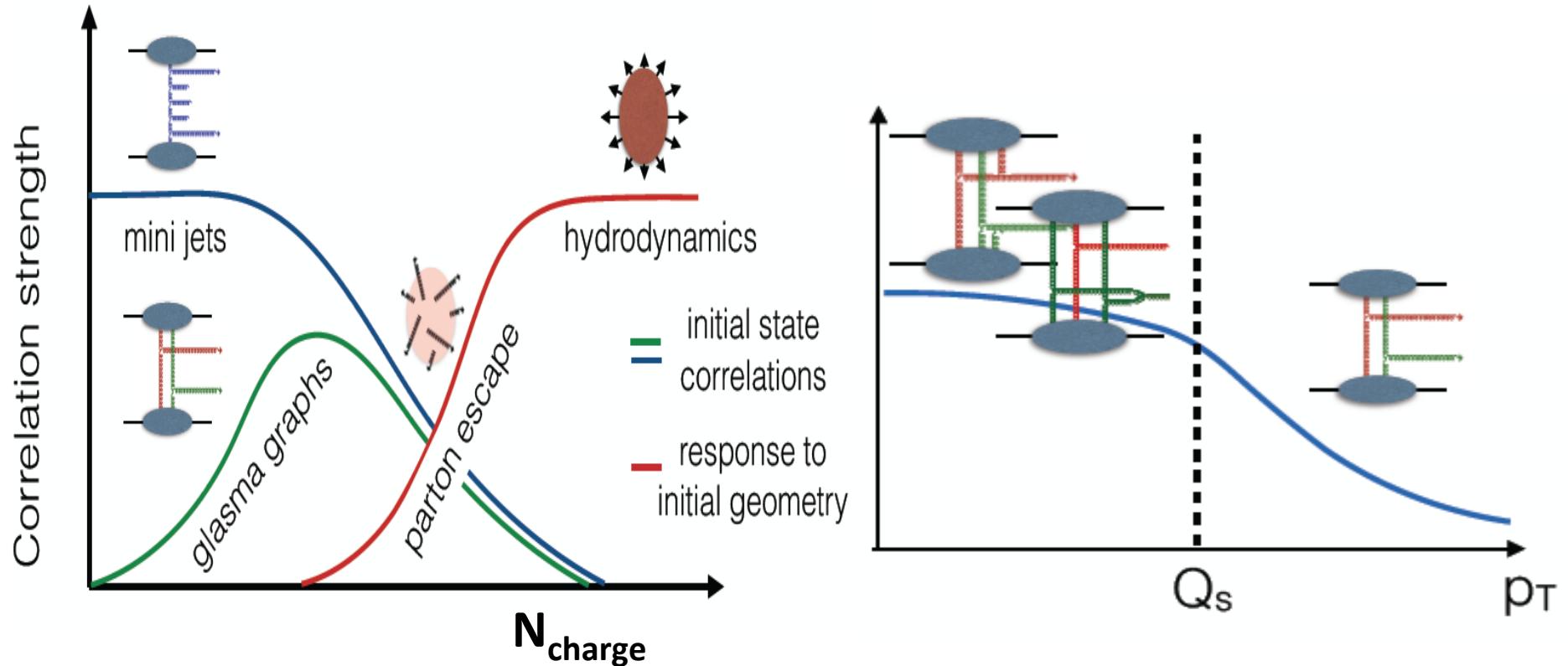
**Hydrodynamic paradigm appears to describe  
multi-particle correlations even in the smallest systems**

**There are however puzzling features of the data,  
questions about the the validity of hydro, fine tuning of initial conditions  
(requiring implicitly strong initial state correlations),  
... and explanation of anisotropies for  $p_T >$  few GeV**

**Initial state QCD frameworks now also able to explain many features of  
the data but systematic treatments are still in their infancy**

**Despite much progress no satisfactory explanation of the data  
-- the problem is still wide open**

## Summary-II



Event engineering across system sizes, energies, and varieties of probes, offers the exciting possibility of exploring dynamical evolution of strongly correlated quark-gluon matter from high occupancy, out of equilibrium, dynamics... to hydrodynamics

Figures: S. Schlichting at Quark Matter 2015

**Thanks for listening!**

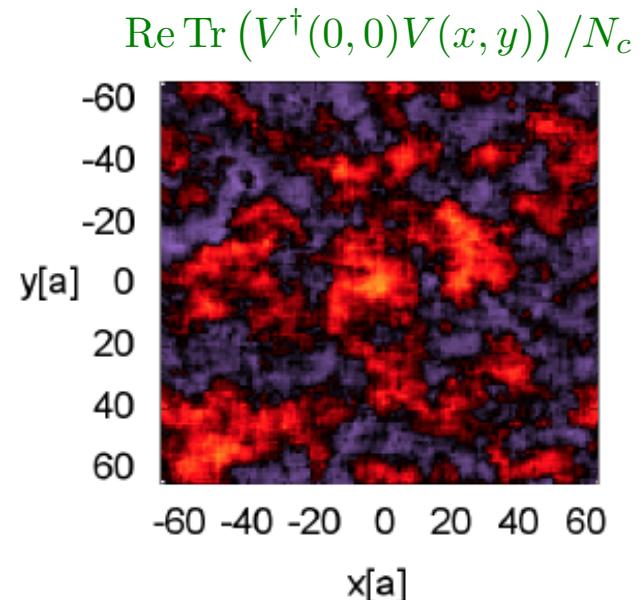
# Higher cumulants from Glasma domains

Dumitru, McLerran, Skokov, 1410.4844

Simple model: express intrinsic higher point correlators as correlators of produced particles with a target field in a color domain, averaged over all orientations of the field.

$$c_2\{2\} = \frac{1}{N_D} \left( \mathcal{A}^2 + \frac{1}{4(N_c^2 - 1)} \right) :$$

$$c_2\{4\} = -\frac{1}{N_D^3} \left( \mathcal{A}^4 - \frac{1}{4(N_c^2 - 1)^3} \right)$$



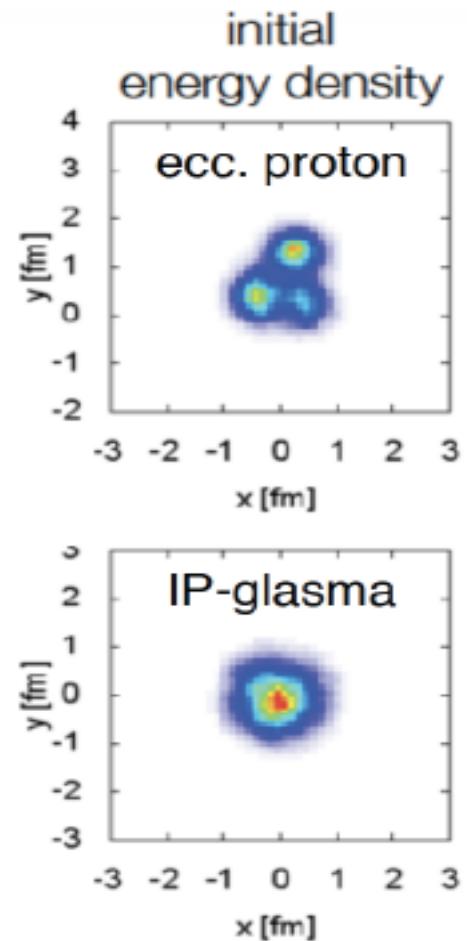
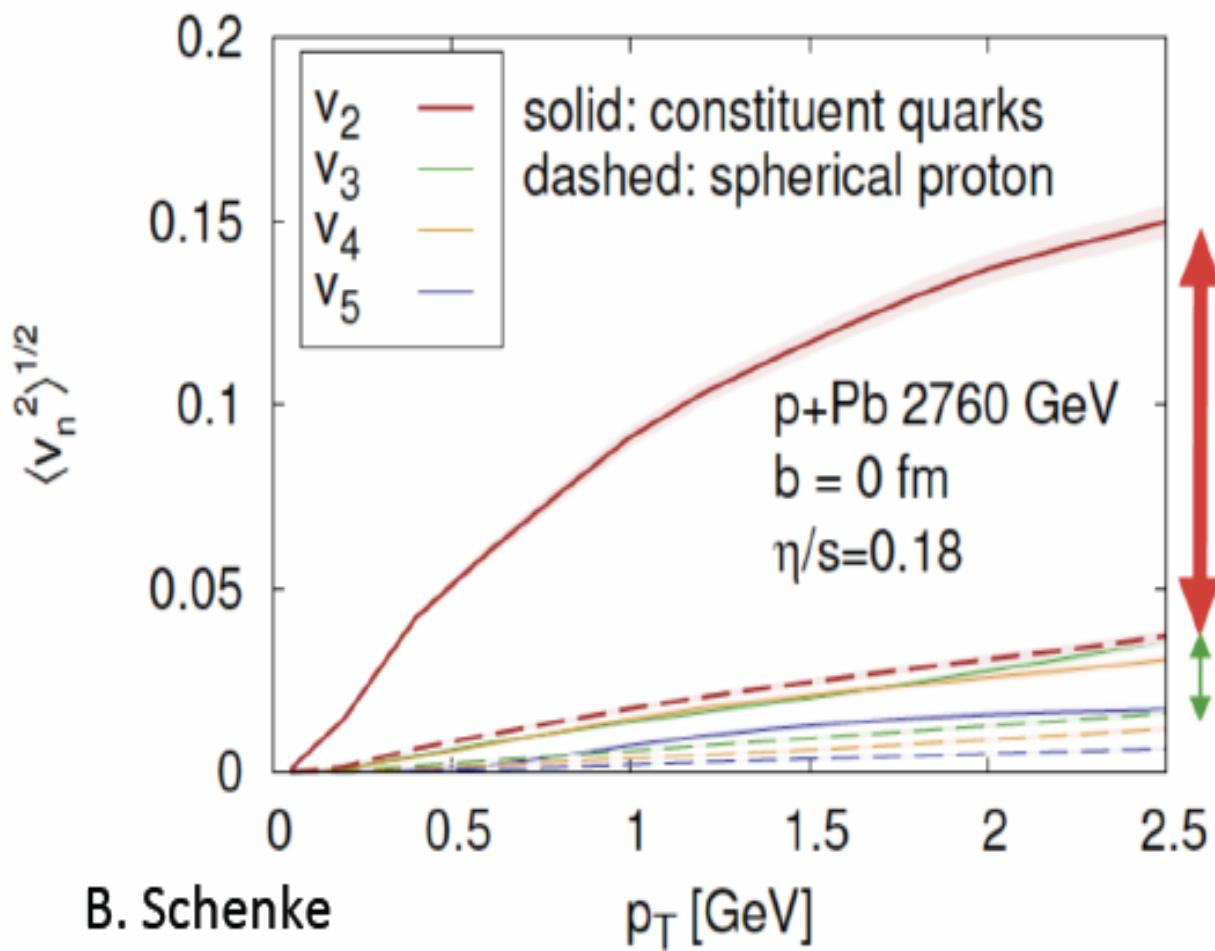
“A” term is the correlation induced between projectile particles due to color field orientation of target (more generically, non-Gaussian correlations)

The  $N_c$  term is the “connected Glasma graph” (Gaussian correlations)

$N_D$  is # of color domains – few in p+A, several in A+A

# Shape matters ?

Schenke,Schlichting, 1407.8458



B. Schenke